Uncertainty quantification

Numerical examples

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Uncertainty Quantification for High Frequency Waves

Gabriela Malenová

Royal Institute of Technology KTH

BIT Circus 2015, Umeå, August 27

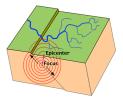
Jointly with Olof Runborg, Mohammad Motamed, Raul Tempone

High	frequency	approximations
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Problem statement

- Propagation of high-frequency waves with uncertain parameters.
- e.g. earthquakes: uncertain medium and source location



Simplified model: scalar wave equation with

- 1. Highly oscillatory initial data.
- 2. Uncertainty (initial data and/or model parameters).

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Cauchy problem for the scalar wave equation

$$egin{aligned} &u^arepsilon_{tt}(t,\mathbf{x})=c(\mathbf{x})^2\Delta u^arepsilon(t,\mathbf{x}),\qquad(t,\mathbf{x})\in\mathbb{R}^+ imes\mathbb{R}^n,\ &u^arepsilon(0,\mathbf{x})=A_0(\mathbf{x})e^{i\Phi_0(\mathbf{x})/arepsilon},\ &u^arepsilon(0,\mathbf{x})=rac{1}{arepsilon}B_0(\mathbf{x})e^{i\Phi_0(\mathbf{x})/arepsilon},\ &t=0,\ \mathbf{x}\in\mathbb{R}^n, \end{aligned}$$

with c wave speed, Φ_0 initial phase, $\varepsilon \ll 1$ wavelength and A_0, B_0 amplitude parameters.

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Sources of uncertainty:

• Wave speed $c = c(\mathbf{x}, \mathbf{y}), \quad \mathbf{y} \in \Gamma \subset \mathbb{R}^N$ random vector,

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with c wave speed, Φ_0 initial phase, $\varepsilon \ll 1$ wavelength and A_0, B_0 amplitude parameters.

Sources of uncertainty:

- Wave speed $c = c(\mathbf{x}, \mathbf{y}), \quad \mathbf{y} \in \Gamma \subset \mathbb{R}^N$ random vector,
- Initial data A₀, B₀, Φ₀.

 \Rightarrow uncertainty in $u^{\varepsilon} = u^{\varepsilon}(t, \mathbf{x}, \mathbf{y}).$

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	Goals	

• Consider quantity of interest

$$\mathcal{Q}^{arepsilon}(\mathbf{y}) = \int_{\mathbb{R}^n} |u^{arepsilon}(\mathcal{T},\mathbf{x},\mathbf{y})|^2 \, \psi(\mathbf{x}) \, d\mathbf{x}, \quad \psi \in C^\infty_c(\mathbb{R}^n).$$

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Proposed method:

- 1. Gaussian beam method.
- 2. Sparse stochastic collocation.

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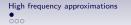
High frequency approximations

Geometrical optics Gaussian beam method

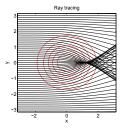
Uncertainty quantification

Stochastic collocation Stochastic regularity for high frequency w

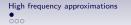
Numerical examples



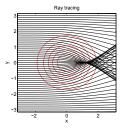
Geometrical optics



- For high frequencies $\varepsilon \ll 1$, computational cost of direct methods grows rapidly.
- Geometrical optics: approximation in the limit $\varepsilon \rightarrow 0$.
- GO breaks down at caustics.
- Remedy: Gaussian beam method.

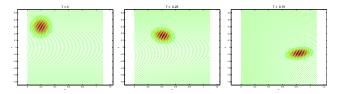


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High frequency approximations	Uncertainty quantification	Numerical examples
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	Gaussian beams	

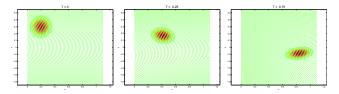


• High-frequency approximation using same Ansatz as GO $v^{\varepsilon}(t, \mathbf{x}, \mathbf{y}) = A(t, \mathbf{x}, \mathbf{y})e^{i\Phi(t, \mathbf{x}, \mathbf{y})/\varepsilon}, \quad \mathbf{y} \in \Gamma \subset \mathbb{R}^{N}$

with Φ and A Taylor expanded locally around GO ray $\mathbf{q}(t, \mathbf{y})$

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- Full solutions on a small $\sim \sqrt{arepsilon}$ neighborhood around the ray.
- Φ has positive imaginary part away from the ray.

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Gaussian beams

First order beams:

$$a(t,\mathbf{x},\mathbf{y}) = a_0(t,\mathbf{y}), \quad \phi(t,\mathbf{x},\mathbf{y}) = \phi_0(t,\mathbf{y}) + \mathbf{x} \cdot \mathbf{p}(t,\mathbf{y}) + \frac{1}{2} \mathbf{x} \cdot M(t,\mathbf{y}) \mathbf{x}.$$

• Require: $\Phi(t, \mathbf{x}, \mathbf{y})$ solves eikonal equation to $O(|\mathbf{x} - \mathbf{q}(t, \mathbf{y})|^3)$ and $A(t, \mathbf{x}, \mathbf{y})$ solves transport equation to $O(|\mathbf{x} - \mathbf{q}(t, \mathbf{y})|)$.

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- We obtain set of ODEs for $\mathbf{q}, \mathbf{p}, \phi_0, M, a_0$.

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- Gaussian shape

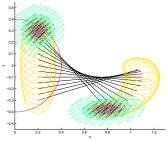
$$egin{aligned} |v^arepsilon(t,\mathbf{x},\mathbf{y})| &= a_0 \exp\left(-\mathrm{Im}(\Phi)/arepsilon
ight) \ &= a_0 \exp\left(-rac{1}{2arepsilon}(\mathbf{x}-\mathbf{q}(t,\mathbf{y}))\cdot\mathrm{Im}(M)(\mathbf{x}-\mathbf{q}(t,\mathbf{y}))
ight). \end{aligned}$$

• $M = M^T$ and Im(M) > 0 for all t > 0 if valid for initial data.

High frequency approximations \circ $\circ \circ \circ \bullet$ Uncertainty quantification

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Gaussian beam superposition



More general solutions \Rightarrow superpositions of Gaussian beams:

$$u_{GB}^{\varepsilon}(t,\mathbf{x},\mathbf{y}) = rac{1}{(2\piarepsilon)^{n/2}}\int_{\mathcal{K}_{0}}v^{arepsilon}(t,\mathbf{x},\mathbf{y};\mathbf{z})d\mathbf{z},$$

 $K_0 \subset \mathbb{R}^n$ compact, $\mathbf{z} \in K_0$ is starting point.

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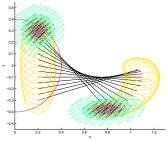
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 $K_0 \subset \mathbb{R}^n$ compact, $\mathbf{z} \in K_0$ is starting point.

- By wave equation linearity, sum of solutions is also a solution
- Accuracy $\|u(t,\cdot) u_{GB}(t,\cdot)\|_E \leq C(t) \varepsilon^{1/2}$

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Stochastic collocation

• Consider quantity of interest

$$\mathcal{Q}^{arepsilon}(\mathbf{y}) = \int_{\mathbb{R}^n} |u^{arepsilon}_{GB}(\mathcal{T},\mathbf{x},\mathbf{y})|^2 \, \psi(\mathbf{x}) \, d\mathbf{x}, \quad \psi \in C^\infty_c(\mathbb{R}^n).$$

where $\mathbf{y} \in \Gamma \subset \mathbb{R}^N$ (random).

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Uncertainty quantification

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• Approximated by

$$\mathbb{E}[\mathcal{Q}^{\varepsilon}(\mathbf{y})] \approx \sum_{k=1}^{\eta} \alpha_k \mathcal{Q}^{\varepsilon}(\mathbf{y}^{(k)}),$$

where α weights associated to the points used.

 NOTE: one full solve of high-frequency problem needed for each y value of Q(y).

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Sparse grid quadrature

$$\mathbb{E}[\mathcal{Q}^{\varepsilon}(\mathbf{y})] \approx \sum_{k=1}^{\eta} \alpha_k \mathcal{Q}^{\varepsilon}(\mathbf{y}^{(k)}),$$

• Key point: choice of collocation point set $\{\mathbf{y}^{(k)}\}_{k=1}^{\eta} \in \Gamma$.

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- Standard quadrature slow when N large.

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- Full grids expensive for $N \gg 1$ (curse of dimensionality).
- Standard quadrature slow when N large.
- Monte-Carlo better but limited to $\eta^{-1/2}$ rate.
- Sparse grids to reduce the cost.
- Sparse stochastic collocation faster if $\mathcal{Q}^{\varepsilon}$ smooth in **y**.

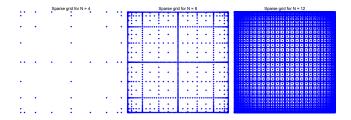
Numerical examples

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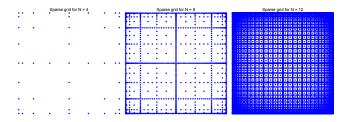
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Sparse grid quadrature



Smolyak sparse grid: nested points on Clenshaw-Curtis abscissas (extrema of Chebyshev polynomials) using total degree index set.

Sparse grid quadrature



Smolyak sparse grid: nested points on Clenshaw-Curtis abscissas (extrema of Chebyshev polynomials) using total degree index set.

- Number of collocation points grows slowly with N.
- Spectral convergence in η (number of collocation points)

$$\mathsf{error} \leq {\mathcal{C}}({\mathcal{p}},{\mathcal{N}}){\mathcal{M}}({\mathcal{Q}}^arepsilon)\,\eta^{-rac{p}{1+\log 2N}}, \quad orall {\mathcal{p}}$$

- Rate depends on smoothness of Q^{ε} : size of $M \sim \left| \frac{d^{\ell}Q^{\varepsilon}}{d\mathbf{y}^{\ell}} \right|$.
- Rate depends only weakly on N.

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Stochastic regularity for high frequency waves

Stochastic Cauchy problem

$$\begin{split} u_{tt}^{\varepsilon}(t,\mathbf{x},\mathbf{y}) &= c(\mathbf{x},\mathbf{y})^{2} \Delta u^{\varepsilon}(t,\mathbf{x},\mathbf{y}), \qquad (t,\mathbf{x},\mathbf{y}) \in \mathbb{R}^{+} \times \mathbb{R}^{n} \times \Gamma, \\ u^{\varepsilon}(0,\mathbf{x},\mathbf{y}) &= A_{0}(\mathbf{x},\mathbf{y})e^{i\Phi_{0}(\mathbf{x},\mathbf{y})/\varepsilon}, \\ u_{t}^{\varepsilon}(0,\mathbf{x},\mathbf{y}) &= \frac{1}{\varepsilon}B_{0}(\mathbf{x},\mathbf{y})e^{i\Phi_{0}(\mathbf{x},\mathbf{y})/\varepsilon}, \qquad t = 0, \ (\mathbf{x},\mathbf{y}) \in \mathbb{R}^{n} \times \Gamma. \\ \mathcal{Q}^{\varepsilon}(\mathbf{y}) &= \int_{\mathbb{R}^{n}} |u_{GB}^{\varepsilon}(T,\mathbf{x},\mathbf{y})|^{2} \psi(\mathbf{x}) d\mathbf{x}, \quad \psi \in C_{c}^{\infty}(\mathbb{R}^{n}). \end{split}$$

- Sources of uncertainty: speed, initial position, wave phase...
- For fast convergence we need

$$\sup_{\mathbf{y}\in\Gamma}\left|\frac{d^{\ell}\mathcal{Q}^{\varepsilon}(\mathbf{y})}{d\mathbf{y}^{\ell}}\right|\leq C_{\ell},\quad\forall\ell\in\mathbb{N}^{N},$$

where C_{ℓ} independent of the wavelength ε .

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Stochastic regularity for high frequency waves

In general, $u^{\varepsilon}(t, \mathbf{x}, \mathbf{y})$ oscillates with period $\sim \varepsilon$ both in \mathbf{x} and \mathbf{y} .

Conjecture/Theorem

The bound

$$\sup_{\boldsymbol{\eta}\in\Gamma}\left|\frac{d^{\ell}\mathcal{Q}^{\varepsilon}(\boldsymbol{y})}{d\boldsymbol{y}^{\ell}}\right|\leq C_{\ell},\quad\forall\ell\in\mathbb{N}^{N},$$

 \textit{C}_{ℓ} independent of the wavelength ε , holds for

$$\mathcal{Q}^arepsilon(\mathbf{y}) = \int_{\mathbb{R}^n} |u^arepsilon_{GB}(\mathcal{T},\mathbf{x},\mathbf{y})|^2 \, \psi(\mathbf{x}) \, d\mathbf{x}, \quad \psi \in C^\infty_c(\mathbb{R}^n).$$

with single family initial data.

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Example 1: Caustics

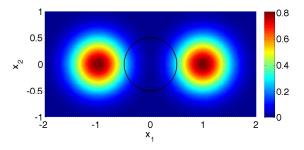


Figure: Two bumps moving towards each other (absolute value).

$$\Phi_0(\mathbf{x}) = |x_1| + x_2^2, \qquad \mathbf{x} = (x_1, x_2).$$

Caustics appear for $t \ge 0.5$. Circle indicates the support of the Qol test function.

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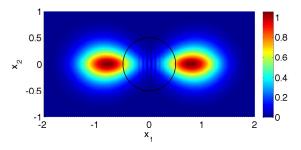


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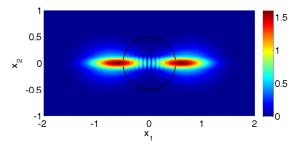


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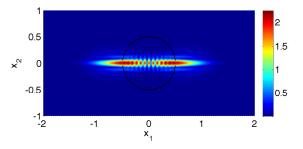


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Example 1: Caustics

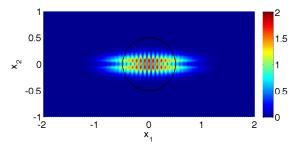


Figure: Two bumps moving towards each other (absolute value).

$$\Phi_0(\mathbf{x}) = |x_1| + x_2^2, \qquad \mathbf{x} = (x_1, x_2).$$

Uncertainty quantification

Numerical examples

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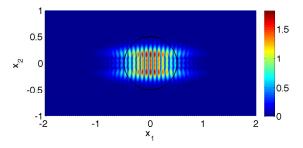


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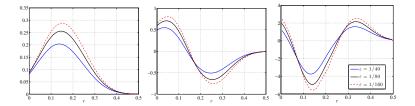


Figure : Quantity of interest $Q^{\varepsilon}(\mathbf{y})$ with its first and second derivatives.

- N = 2 random variables initial position (y_1) and constant speed (y_2) .
- Qol along the line y(r) = (1 + r, 1 + 2r), for different wave lengths ε.

Uncertainty quantification

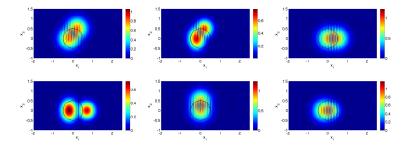
Numerical examples

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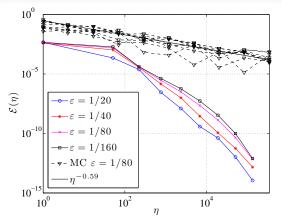
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Example 2: Sparse grids



- *N* = 5 random variables (speed, initial data pulse shape, position).
- $\Phi_0 = |x_1|$.

Uncertainty quantification



Relative error in the expected value of $\mathcal{Q}^{\varepsilon}$ for levels $\ell \geq 1$:

$$\mathcal{E}(\eta(\ell)) := \frac{\left| \mathbb{E}[\mathcal{S}_{\mathcal{I}(\ell_{\mathsf{ref}})}[\mathcal{Q}^{\varepsilon}]] - \mathbb{E}[\mathcal{S}_{\mathcal{I}(\ell)}[\mathcal{Q}^{\varepsilon}]] \right|}{\left| \mathbb{E}[\mathcal{S}_{\mathcal{I}(\ell_{\mathsf{ref}})}[\mathcal{Q}^{\varepsilon}]] \right|}.$$

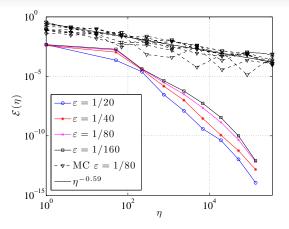
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Uncertainty quantification

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990



- Fast spectral convergence compared to Monte-Carlo.
- As ε decreases, error converges \Rightarrow uniform bounds.

Uncertainty quantification

Numerical examples

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THANK YOU FOR YOUR ATTENTION